

today we shall start again recapitulating
 what we do it in the last lecture the new things those we learnt and then we
 will
 do something more with kinetic theory and plan is that towards the end of
 today's lecture hour i
 will tell you briefly about non-ideal situations why non ideal situations are
 important it is
 important because of the fact that you know ideal gas what we are dealing
 so far there is
 no interaction and if there is no interaction there is no phase transition
 possible most of
 the situations only one exception to be correct i tell you that without
 interaction phase transition
 happens that is called bose einstein condensation i quoted this name because
 it involves the name
 of a great indian scientist satyanranath bhosh but that is essentially a
 quantum phenomena
 happening at very low temperature where there is no interaction in that sense
 it is the
 nature of both particles or rather to be precise the statistics of both
 particles that give us
 the phase transition
 so called formation of bose condensate which at least you know about
 them in popular level at least but what i want to tell you at the moment that
 we are dealing with ideal gases and these ideal gases cannot
 lead to any phase transition because there is no interaction
 so first and
 foremost criteria of a phase transition should be interactions between the
 particles i am dealing
 with
 so that is why i will bring in some non ideal nature and tell you what is
 vapor and how does the
 pv diagram for a liquid gas transition looks like but before that lets start
 with what i wrote
 here recapitulating what we learned new thing i talked about pressure ok
 pressure is
 related to $m n l \text{ cube } v r m s \text{ square}$ i remind you what is $v r m s \text{ square}$ v
 $r m s \text{ square}$ actually it is a vector dot product if you like i can
 decompose it in three components $z \text{ square}$ and divided by one over n
 so that its the
 average ok its the average and i am summing over all the particles i have ok
 so this is the
 average information that is carried here and this is mass of each gas molecule
 n is
 the number of particles in the container and $l \text{ cube}$ is the volume of the
 container
 where i chose a cube but it is not essential to choose a cube you can have a
 sphere if
 you like and this you can write further as p is equal to one third $\rho v r m s s$
 square or
 i can go further $p v$ is equal to one third $m v r m s s \text{ square}$ which i can write
 also as if
 my mass of each molecule is m $m n v r m s \text{ square}$
 so this is i will drop this

expression v_{rms} in the subscript from the time being and I will be assuming that I am dealing with average velocities which are root mean square velocities now question was how do you relate it to temperature so have pressure volume but they are related to microscopic objects which we can never measure in experiments what we measure with a thermometer is temperature so it is important that this relation should be connected to temperature so what I did first ideal gas equation of state what is an equation of state equation of state is something that connects the different thermodynamic variables or I am talking about chemical system involving p , v and t so my equation of state will be something that connects pressure volume and temperature this is true for an ideal gas and this is a combination of Charles law which says either p is proportional to t if I keep volume constant or v is proportional to t if I keep pressure constant or its Boyle's law which tells me that pV itself is a constant if I keep temperature a constant ok these can be verified for a gas which is at very high temperature and very low density which is the limiting situation for any real gas where I can approximately treat it as an ideal gas ok so now to I put these things together this equation I got here this is my ideal gas equation of state $pV = nRT$ I can simply write $pV = n k_B T$ and then what we have found out very important expression that $\frac{1}{2} v_{rms}^2$ is equal to $\frac{3}{2} k_B T$ and then this gives me very important relation that temperature is defined in terms of average kinetic energy of the gas molecules this I can physically interpret in the following way if I increase the temperature I expect the gas molecules to have more kinetic energy so if you can think of this way $k_B T$ has dimension of energy these I will call the thermal energy if you have higher and higher temperature they will have more kinetic energy the gases will move with more kinetic energy this is something I physically expect and this is what kinetic theory tells me about so once I have this I have a new definition of temperature so far we were dealing with definition of temperature which we measure in thermometer that is in our calorimetry but kinetic theory gives me a more fundamental definition of temperature that mean translational kinetic energy is related to $\frac{3}{2} k_B T$ ok now these two taken together $pV = nRT$ and kinetic energy for n molecules if I have n molecules total kinetic energy will be $\frac{3}{2} n k_B T$ if I have both of them together

then I arrive at an expression pV is equal to two thirds of the translational and its all kinetic I told you why it is kinetic its kinetic because I am assuming ideal gas and ideal gas does not have any potential energy because there is no interaction between the molecules this is the first point and second point I said that it is all translational because I am dealing with a monoatomic gas and there is only translational degrees of freedom ok so if I increase the temperature translational degrees of freedom translational energy increases and that defines my temperature this is an equation in which I have completely gotten rid of I have completely gotten rid of temperature I am connecting pressure volume to the average kinetic energy of the system ok and secondly I mentioned briefly that this is a very sacrosanct equation because pV is equal to $n k T$ is valid for ideal gases which are classical which are very high temperature low density situation I cannot write a similar equation if I go to very low temperature but as long as I am dealing with a monoatomic gas and as long as I am treating ideal systems that means there is no interaction I arrive at this equation even the Bose gas I was mentioning before you can have an equation like this and I told you that this two and three these two comes from the fact if E_p energy of a particle simply I am doing Newtonian mechanics so any particle with momentum p with energy p^2 by two m so there is a two factor coming here and three coming in because of the factor that I have considered three components of velocity usually our world is in our conventional physics classical physics world is three dimensional as three components of velocities and there is no difference between these three components of velocities there is something called isotropy which I may refer to it many times now I told you very something very very important since these three are completely equivalent then I can have an equi partition of energy ok I can have an equi partition of energy ok what do I mean by equi partition of energy equi partition what do we mean by equi partition of energy that my v^2 I can think of its consisting of three pieces v_x^2 and v_y^2 and v_z^2 this is my v^2 ok so kinetic energy has contribution from x component of velocity y component of velocity and z component of velocity so since v_x^2 v_y^2 v_z^2 are equivalent in the probabilistic sense these are all rms I am talking about I can write it as

three v_x^2 in the average sense ok since i can write it as $3 v_x^2$ i get $3 v_x^2$ if i multiply by $\frac{1}{2} m$ is equal to $3 \times \frac{1}{2} m v_x^2$ so k always stands for the Boltzmann constant even if sometimes i may forget to put this k_B subscript here though immediately give me $\frac{1}{2} m v_x^2$ is equal to $\frac{1}{2} k_B T$ so you see the kinetic energy if you like contribution to the kinetic energy for the x component of velocity $\frac{1}{2} m v_x^2$ is equal to $\frac{1}{2} k_B T$ similarly there is nothing great about x i can write $\frac{1}{2} m v_y^2$ as $\frac{1}{2} k_B T$ this is equipartition which i said that three components of velocity three components of velocities ok and energy is apparently $\frac{1}{2} k_B T$ for each molecule on an average in each direction ok well what i mean by each direction is clear that $\frac{1}{2} m v_x^2$ is equal to $\frac{1}{2} k_B T$ $\frac{1}{2} m v_y^2$ is equal to $\frac{1}{2} k_B T$ and these three adds and gives me $3 \times \frac{1}{2} k_B T$ if i have for n molecules then $\frac{1}{2} m v_x^2$ ok will be first $\frac{1}{2} n m v_x^2$ translational kinetic energy associated with v_x this will be this quantity which will be $n \times \frac{1}{2} k_B T$ ok so this is called the equipartition energy whenever you have this type of from ideal gas then you can ask me that what is the energy associated with each degrees of freedom i should be very very careful when i mean mean what do i mean by degrees of freedom ok let us recall a bit of classical mechanics ok in classical mechanics if a particle is moving along a line let us say x axis then i can immediately characterize the particle by x and also velocity v_x ok if it is under a force then i can say x and v_x so x if you like is the degrees of freedom ok you can also think of v_x is one of the degrees of freedom if you have an x if you have a v_x so there is one degree of freedom ok now if particle is allowed to move on the floor of this room you have a v_x similarly you have a v_y right so i will say x, y that gives me the coordinates it has two degrees of freedom similarly if i have a particle moving anywhere in this room i will need three coordinates x, y and z and correspondingly three components of velocities v_x, v_y and v_z so essentially i have three degrees of freedom and i am saying for each particle on an average ok on an average if temperature is T i am talking about equilibrium situation if temperature is T for each particle on an average for each degree of freedom that means associated with each v_x here x has no role to play because its free particles ideal gas

for each
 v_x will have an energy average energy which is $\frac{1}{2} kT$ this is called equi
 partition
 of energy and this is very very important now if you think of n such particles
 where
 this is your capital and you know you will need many many three for three for
 each three
 for each particles these I explained for one particle now you may have n
 particles so
 n particles each moving in three dimensions
 so I will have v_x v_y v_z ok v_x v_y and v_z and now for
 each of
 this average kinetic energy I will have $\frac{1}{2} kT$ ok n number of particles because
 they are not
 interacting and I will have same term for v_y $n kT$ and I will have a similar
 term $\frac{1}{2} n k$
 bT for this and the total will be three by two $n kT$ ok this is how I get
 to this expression using the equi partition or rather kinetic
 theory tells me about the equi partition theorem now I can go further I told
 you there
 could be other special cases for example I can get harmonic oscillators which
 have energy $\frac{p^2}{2m}$ and this is the potential energy ok this is the
 form of potential energy for a
 harmonic oscillator and if I have this potential energy then energy of this
 average energy of this
 harmonic oscillator when this oscillator is at a temperature T is kT $kB T$
 this is equi partition
 applicable to harmonic oscillators ok whenever you have this quadratic form ok
 here ideal gas I
 was only having translational degrees of freedom but here you see I have x and
 p both both I can
 treat as independent degrees of freedom both are contributing to my energy and
 both is
 quadratic in form this is very very important
 so so for ideal gas I was calling my
 degrees of freedom only as v_x v_y v_z ok but when it comes to harmonic
 oscillator
 where force is of the form $-kx$ and potential is the of the form $\frac{1}{2} kx^2$
 square
 as we have learnt in my mechanics course then total energy is this form and I
 have
 effectively two degrees of freedom which are p and x and both giving me
 $\frac{1}{2} kT$
 giving me average energy of an harmonic oscillator is $kB T$ ideal gas
 molecules it is $\frac{1}{2} kBT$ and
 harmonic oscillator it is kBT I just argued it out using the equi partition
 which I try to
 invoke upon you using the kinetic theory of gas now this is the average for
 harmonic
 oscillator if you are in non interacting this word is very important non
 interacting
 harmonic oscillators then simply you can write this now you see these harmonic
 oscillators
 are not oscillating in one dimension rather in three dimensions then you will

have a term
 three sitting here ok that tells us a lot about equi partition of thing i
 explain to you
 what each degrees of freedom means and then i will proceed further to
 calculate something
 measurable in the experiment which i briefly touched upon in my last lecture
 that ok i talked about specific heat
 and specific heat can be of two types firstly you calculate ah specific heat
 keeping volume of the container fixed that i call c_v ok this is c_v ok and
 then i can
 have it keeping the pressure constant measure the specific heat which i can
 write
 in general c is equal to $d e$ sorry $d e d t$ ok you know heat energy specificate
 you define
 in calorimetry in terms of amount of heat that is being absorbed or released
 in raising the
 temperature here we know heat is an energy and we are dealing with energy here
 specific heat is
 equivalently related to change in this derivative simply tells me this is the
 change in energy
 average energy if i change temperature or if i use more familiar notation for
 you it is $\frac{d e}{d t}$
 t ok you change temperature by an amount Δt and that this is the change
 in energy you take
 the limit Δt tending to zero you can write in your calculus form if you
 keep volume fixed
 then it is called c_v if you keep pressure fixed it is called c_p and for an
 ideal gas i just quote
 it without any proof or further for ideal gas $c_p - c_v = R$ this i will try
 to
 elaborate more and there is a ratio γ which we will calculate c_p by c_v
 which i will
 use when i go into thermodynamics i talk about adiabatic processes then these
 things will be
 useful now if you do this you can immediately see c_v let us focus on only c_v
 for an ideal gas it should be three by two $n k_B$ which is
 called the dulong petite law if you do it for an harmonic oscillator using the equi
 partition you do it for an harmonic oscillator this expression three $n k_B$ which
 i have
 wrote here gives me c_v is equal to three k_B t three $n k_B$ ok
 so you see these are called
 dulong petits law ideal gas it is three by two i told you in the last class
 this three comes from
 dimensionality of the thing and here it is three because where do you get the
 other half factor
 form in the harmonic oscillator because in this harmonic oscillator if there
 is a half $k x^2$
 which is very very important because now you have two degrees of freedom in
 that sense one given by
 x and one given by p both contributing half kt to my average energy
 so harmonic oscillator this is
 c_p this is c_v and you can calculate $c_p - c_v = R$ using this relationship you have
 here you remember R

is equal to $n k_B$ but coming back to the ideal gas
 so when i talk about cpv its ideal gas
 harmonic oscillator it is just three $n k_B$ now coming back to ideal gas
 so far
 we have been talking about mono atomic and hence i was talking about only
 translational degrees of freedom and now if i have a diatomic cache ok
 so go beyond
 mono atomic you do diatomic cache okay all you have to do remember how many
 how many degrees of freedom do you have and how does this degrees of
 freedom contribute to your energy degrees of freedom contribute all degrees
 of
 freedom contribute half kT to energy ok
 so diatomic case diatomic
 so first thing is translational
 again this molecule is moving in a three dimensional container
 so i know again
 there is no interaction
 so its ideal gas freedom container and i do not consider any interaction so
 one is translational degrees of freedom how many three
 so what do they contribute three
 by two $k_B T$ ok but thats not the end of the story ok let us take an oxygen O_2
 ok
 so you have
 two atoms ok there are translational degrees of freedom but of course if i
 take a dumbbell
 like shape i know molecules can rotate if i take an axis it can rotate about
 this axis
 i know what is rotational kinetic energy i have done it in my rotation chapter
 in
 mechanics
 so i have a rotational kinetic energy and what is this these are two axis of
 rotation if
 i rotate about one two possible rotations and this gives me half $I \omega^2$
 square and half $I \omega^2$
 two square
 so i have fixed this axis and i am allowing rotations i can have 2 rotations
 so those
 both of them giving me half $I \omega^2$
 so i have 2 degrees of freedom here and remember this
 is also what i say quadratic in nature if that be the case i immediately know
 that i am going
 to get a three by two kT here i am going to get half kT for each rotational
 degrees of freedom so
 total i will be having total average energy for a diatomic molecule is f by
 two $k_B T$ and then
 specific head if i take n such molecules average energy will be five by two n
 $k_B T$ giving me
 a specific head that is $\frac{5}{2} n k_B T$
 so you see that number of degrees of freedom i increase i get
 a different expression for my total average energy this is the total average
 energy
 for n such diatomic molecules and once i put that in the expression of
 specific heat which i experimentally can measure ok that goes into five by two
 instead of three χ

two next question you will ask what will happen as a polyatomic molecules
 so for a polyatomic molecule you have to be bit more careful and one has to go further than this diameter now
 let us go to polyatomic now i have many many atoms forming a molecule polyatomic situations
 and i will use this equi partition theorem which i talked about roughly if you have a polyatomic molecules you can think of its as a rigid body for example firstly crudest possible approximation let us consider its at a rigid body ok now if it is a rigid body you know rigid body mechanics teaches us a rigid body has six degrees of freedom ok six degrees of freedom y six you may be knowing only a rigid body rotating about a fixed axis and you know l is equal to $i \omega$ but truly speaking a rigid body can rotate about a fixed point what do i mean by that its rotating about a fixed point means instantaneous axis of rotation passes through this points ok or since that can happen at the same time i can always give a translation to any particular point you know the concept of center of mass so let us focus on centre of mass i can say rigid body is rotating about its center of mass and then the center of mass may have a translational degrees of freedom and it can move in any direction so there will be three translation and three rotations okay since you may not be very familiar with this 6 degrees out of freedom business let us invoke upon you that how does this 6 come ok what is the definition of a rigid body definition of rigid body tells you that distance between the any two points in the rigid body is fixed ok any two points is fixed so if i put particular one particular point a mass here then it has three degrees of freedom because it has no constraint with it this is the my first particle this is my second particle this is my third particle ok now first particle has three degrees of freedom three what about the second particle second particle can do whatever it likes but it has to be at a fixed distance from the first particle so second particle has degrees of freedom too ok because it has to maintain a distance as i said in the diatomic molecule case ok two two and now third particle has to keep a fixed distance from two and one both so it can move only in three dimensions actually but always keeping this constraint that it has to be at a fixed

distance from two and from one
 so it has one ok this gives me six total degrees of freedom of this
 system of particles but if i what about if i bring a fourth particle here the
 fourth particle ok
 well fourth particle has to keep a fixed distance with all the other particles
 in the business
 it has to keep a fixed distance from one two and three
 so it does not have any free degrees
 of freedom in that sense
 so total degrees of freedom is six i told you that 10 to the power
 23 particles if you have to deal with you cannot write 10 to the power 23
 second order differential
 equations with given the force form and you cannot solve it for that matter ok
 so but what helps us
 in mechanics is this rigid body approximation so this is a rigid body
 approximation which you
 may not be very familiar with that is why i am spending some time in in this
 rigid body thing
 rigid body approximation why it is approximation there is no ideal rigid body
 in world you already
 know i am pretty sure of that einstein's theory of relativity tells us that no
 information can
 propagate at a speed which is greater than speed of light
 so i cannot have anything instantaneous
 in our real world ok there is always a time required but rigid body assumes
 that if i give
 a disturbance to any point of rigid body ok any point of the rigid body then
 rigid body should
 information should propagate to other part of the rigid body instantaneously
 which is not possible
 and that is why rigid body is an approximation but that is a very good
 approximation and
 secondly i have 10 to the power 23 particles let us say in a rigid body but i
 deal with just 6
 degrees of freedom three translational and three rotational and you know that
 simplifies life to
 a great extent ok if that be the case what is the polyatomic molecule energy
 there are six
 degrees of freedom each giving me half kt and n of such molecules i will
 simply have three
 $n k t$
 so number of degrees of freedom is actually doubled ok
 so specific it will be $3 n k$ but that
 is not the end of the story i am assuming i say its a rigid body approximation
 there could be some
 vibrational mode and if you count this vibrational modes there could be f such
 vibrational modes
 so you will be having for each vibrational mode you will have three kt that is
 coming from my
 rigid body thing and then i will have f by two $k t$ for possible vibrational
 mode of the rigid bodies
 and this will be my total energy and accordingly this specific heat gets
 modified
 so this is

all i wanted to discuss about specificate if we partition and if you follow your book you will see there are many discussion on this and also you can refer to professor a c vermont's book where also things are done at some length so what i have done so far that p is equal to again i write $m n v^2$ square this is $v r m s$ square and kinetic energy gives me the proportionality to temperature v^2 square is proportional to $k v t$ if you like ok this is a very fundamental thing of ah kinetic theory of gas what i am dealing with now if i have a different approach which actually i pick up from professor harish verma's book you can see if you have one gash for which $p v$ is equal to $m n v^2$ square now you can motivate yourself that actually temperature will be a function of half $m v^2$ square ok this functional form i derived using the ideal gas equation ok i use the ideal gas equation to arrive at this functional form physically i argued that temperature is increasing mean kinetic energy of the molecules i am considering they must also increase and this is the what this functional form tells us so it should be a proportional from my physical arguments but this can be elaborated in the following way also if i take the same gas in the same container at a different temperature and pressure which i know for example this temperature which i measure could be if you like some temperature which is of the order of two seventy three point one six absolutes and you can write $m n v$ naught square ok then you immediately know p by $p \theta$ will be equal to v^2 square i remind you these are all rms's rms velocities so v^2 square by v naught square once you have this form now you recall your charles law p is proportional to t for a given volume you can write v^2 square is equal to p by p naught into v naught square and if you remember your charles law and if you know that p naught and v naught are constant for a gash ok you immediately arrive at something which is v^2 square is equal to v naught square by t naught into t so that gives you precisely the proportionality i am talking about and then kinetic energy if you v^2 square if you extract the kinetic energy from this this is a constant for a gas given gas this is a constant so you get v^2 square is proportional to t and now hydro is said this proportionality just by argument hand working argument if i make it energy and dimensionally this quantity should be bohemian constant and i can have some number this number i are good to be three by two ok because dimensional analysis argument cannot give me

whatever this number
 a is ok having done all these things
 so far i will proceed and tell you little bit about
 ah how to use kinetic theory for different situation and that will probably
 end today's
 whatever i wanted to do
 so if i mix two gases two gases but everything is in
 equilibrium i call it thermal equilibrium which already tells me there are
 other equilibriums and it will be discussed in the thermodynamics
 part of this lecture series thermal equilibrium means temperature everywhere
 in the gas is the same
 so i do not need to bother about the temperature which is same everywhere
 and if it is this i have two gas molecules one is having mass m one other is
 mass $2m$ i must have
 $m_1 v_1^2$ again in average sense that's $\frac{1}{2} m \overline{v^2}$ ok
 so this is an important
 thing if those two gas molecules are in equilibrium i must have this condition
 satisfied
 on an average ok from this you can calculate the rms velocity if i give you
 the temperature and
 rms velocity of one of the gas molecules you know what you can mathematically
 calculate
 what is the rms velocity of the other one ok
 so now second will proceed further ok
 now let us proceed do i have boyle's law $p v$ is equal to constant negative
 temperature what did i find out for $p v$ ok $p v$ is related to as i wrote one third
 capital m one third capital m small $n v^2$
 square average ok now this is my $p v$ now i have already told you this fellow
 is nothing but proportional to temperature ok we have already used boyle's law
 in getting
 into this equation p is equal to one third m and v^2 but i will make
 thing slightly
 more complicated ah rather i am doing a self consistency check here just i
 am saying that if this quantity is temperature proportional to temperature
 so this
 right hand side must be equal to a constant if temperature is constant
 so $p v$ is constant if t is constant because right hand side of this
 equation is entirely proportional to temperature and this gives me my boyle's law
 similarly you can
 argue about charles law if you keep p fixed ok and this quantity is
 proportional to temperature
 so if you keep v fixed ok
 so p is proportional to temperature that is your charles law ok now
 let me see what else we can talk about okay in ok now have i got this law what
 does avogadro say that given t that means i fixing the temperature i am p equal
 volumes of all gases have equal number of particles that's what we say the
 avogadro number
 if we call one mole of the sample
 so equal volume given temperature given pressure
 i am fixing temperature i am fixing pressure ok now what is $p v$ equation for
 the first
 $\frac{1}{3} n m \overline{v^2}$ ok well sometimes i will put this bar to
 remind you
 that these velocities we are talking about are the average velocities now for

the second gas i
 do have $p_i v_i$ is equal to one third $n_i k_B T$ $m_i v_i^2$ square this is my $p_i v_i$
 equation ok volume is
 same but there is one more to this which i wrote earlier this only tells us
 nothing
 because i have quantities involving n_i one $m_i v_i^2$ to arrive at this
 claim which avogadro made ok avogadro's hypothesis that if i want to do that i
 need
 one more since temperature is fixed T is fixed i must have this quantity which i
 wrote few minutes back that
 when you have two gashes at same temperature this must be satisfied once you
 have both these
 together put both these equations together which gives me n_i is equal to n_i
 two
 so completely see charles and boyle's law were somehow expected to come out of
 it because
 in arriving at this equation at some point $p_i v_i$ equation some point i used
 ideal gas ok to at
 least get the energy p_i is equal to one third $m_i n_i c_i^2$ square $p_i v_i$ is equal to one
 third $m_i n_i c_i^2$ square or
 $m_i n_i v_i^2$ square whatever i derived from the kinetic consideration but to connect
 it to energy
 $p_i v_i$ is equal to two third e_i what i needed is ideal gas equation or some way i
 inspired
 that using the definition of absolute scale and the thing i did from professor
 a c
 verma's book but what is not expected here ok that it also gives me the
 avogadro
 hypothesis whatever other laws of thermodynamics laws of thermal physics i
 know can be arrived at
 similar fashion ok
 so you see that microscopic physics what i am doing here ok are all will
 follow or lead to the macroscopic equations i know from my school early school
 days for
 example second examples will be for example ah dalton's law of partial pressure
 dalton's law of partial pressure which says that you have many types of gases
 let us assume all of them mono atomic and then i put them in a container ok
 so i put them in a
 container and ask what is the pressure pressure in the container ok its called
 partial pressure
 law which tells you that if i have for example some types of gas molecules
 then this total
 pressure which is p which is being exerted on the wall of the container have
 many pieces
 p_1 p_2 p_3 corresponds to fast cache fuzzy physically what it means
 that if i only
 had ok the gash number 1 in the same container the pressure exerted by the gas
 number 1 to the
 container ok would have been ok p_1 similarly if there are no one two one
 three etcetera only
 second type of gas then the pressure exerted will be p_2 but when all of
 them are put together
 then its p_1 plus p_2 plus p_3 and
 so on well this is partial pressure you should bear clearly

in mind what i mean when i say p_1 or when i say p_2 if i remember how do we define pressure pressure we defined in terms of momentum transferred right momentum transferred per unit time so momentum transferred per unit time divided by l^2 square this area of the container so it was important that momentum transferred ok now first molecules they are all independent non interacting goes and hits this wall ok that will lead to a momentum transfer Δp_1 and if i take all of this first molecules then the momentum transfer will be simply f_1 ok now second one will again give me a momentum transfer of f_2 so net momentum transfer will be summation over i f_i this is the net momentum transfer because of the all molecules present in the container per unit time and net momentum transfer will be this quantity i have written with summation here where some extends over all the molecules i have and then if i calculate pressure which is just dividing the net momentum transfer per unit time which i was interested in so pressure that this system of gas molecules with excerpt on the wall of the container will be given by net momentum transfer by l^2 square and which you can see see p_1 p_2 etcetera so p_1 is just due to momentum transferred by the gas molecules one p_2 is just the momentum transfer by the gas molecule two and so on third thing or the last thing in this regard so that tells us the physical origin of the dalton's law of partial pressure finally if two gashes two gases maintaining same pressure and temperature we allow it to diffuse allow to diffuse you can very easily argue that diffusion rate what should be the diffusion rate how fast they diffuse it diffuses because they have a rms velocity and they want to spread out so one molecule spreads into the other now rate if i say can simply be proportional rate simply be proportional to r_1 by r_2 simply be proportional to v_1 again rms remember whatever i am talking about rms and we already have seen rms $m v_{rms}$ is simply given by $3p$ by ρ so if you substitute this the rate of diffusion will be inversely proportional to density so how do you get it i repeat the argument the argument is the following that two gases same pressure and temperature are allowed to diffuse one is diffusing into the other gas is coming out of a cooking gas cylinder and diffusing into air so i expect the rate of diffusion how fast it diffuses ok now r_1 by r_2 should be proportional to how fast the gas molecules are moving

so that will be given in terms of these rms speeds and if this is given in terms of rms we know the expression of rms okay if i substitute it back here i will find out this rate is proportional to $\rho^{-1/2}$ inversely proportional to the square root of density ok this is called the grams law of diffusion all these things we have say are experimentally measured ok and kinetic theory though in kinetic theory we are talking about a microscopic description microscopic description given in terms of velocity distribution which i told you that there is a speed distribution there is a velocity distribution there is an rms speed but all the experimental observations which we make can be reflected here since i talked about last two minutes of today's thing since i have talked about so much about the diffusion in last two three minutes i will try to tell you something now we have assumed all through which i will complete in the next lecture that i have been telling you that see i am dealing with gases which are not colliding except for the wall when i calculated if you remember the momentum transferred and hence the pressure i was specific this particle is not having any collisions it is successively bouncing back and forth between the two walls of the container and this is obviously a very very drastic assumption ok now there is a concept if i keep a gas here there is a concept its called mean free path ok mean free path cache molecules collide and there is a concept of mean free path if you see gas cylinders gas leaking out of gas cylinders they do not go uniformly ok if you could see ok by a microscope how the molecules are moving they will have some zigzag motion because there are collisions between the molecules between the molecules there are collision between the gas molecules this is an important thing what is mean free path mean free path is average again the word average is very important average distance that a gas molecule traverses between two successive collisions two successive position this is a very important concept because you know what we have been doing is very idealized we have to come closer to real world and you can find out that which i will do in details in the next lecture that its roughly proportional to $\frac{1}{n \pi d^2}$ ok d square with some constant here which i cannot find out ok you can check whether i am writing this expression correct density of particles density of particles and this is the diameter of the molecules so i am assuming molecules have a finite size ok assuming molecules having a

finite size

this is the diameter of the molecule this is the density of molecules now if you assume the limit this λ will prove in the next class I am just telling you a bit about it if we assume the limit n goes to zero or d is zero ok they are very very small impossibly compared to the other length scales of the problem this quantity is very very large and the approximation I was making that there is no collision ok between it hits two walls ok and I neglected all the collisions that is not true if n is very large and d is appreciably large if they are very small this fellow goes very half very high and I can make this approximation so here I stop telling you that there is a role of finite size ok and you see I always talked about low density so that mean free path is very very large ok there is a size of the there is a size of the gas molecules that has to be taken in consideration that brings in the concept of mean free path that is useful in many many cases even if you consider free electrons in a metal and you consider the conductivity you use some free electron theory and that free electron theory also needs a concept of mean free path I will stop it here today thank you