

ok welcome to the second lecture of thermodynamics in which as usual we shall start with a recapitulation so we discussed thermodynamic variables which are macroscopic objects measurable in experiments there are of two types one is extensive and other intensive extensive thermodynamic variables are indicative of the system size for example number of particles volume internal energy later we will introduce probably concept of free energy entropy these are extensive variables they are indicative of the system size if we double the size of the system maintaining equilibrium these quantities will be doubled on the other hand pressure temperature these are intensive variables these quantities do not depend on the system size okay insensitive to the system size they are fixed by the equilibrium condition similarly if I have density which is n/v this quantity is an intensive quantity even though n is extensive and v is extensive so these extensive and intensive variables are my thermodynamic variables which describe a system but the system which I am considering is always interacting with a reservoir or we called it the rest of the universe or simply the universe so system separated from the universe by walls we talked about different type of walls for example adiabatic wall which does not allow any heat exchange or diathermic wall where I can have heat exchange in general I can have both heat exchange and mechanical interaction so system universe and system is separated from the universe by wall system is small whereas reservoir is very big I can assume that heat capacity is infinite which I mean by saying that if I extract heat from this or wire or I release some heat to the reservoir its temperature does not change ok so as I said walls are diathermic and adiabatic in general we will consider walls which are diathermic and also movable so that I can do some mechanical work on the system on the other hand system can also do some mechanical work on its own on the universe so do two types of interactions I mentioned mechanical and thermal there could be other type of interactions for example there could be particle exchange which we are not interested in in the present set of lectures then the concept of equilibrium ok equilibrium means all the thermodynamic variables we measure do not depend on time this is an idealized concept all these walls there is no

perfect adiabatic wall in this universe ok this is an idealized concept
similarly equilibrium is an idealized concept ok i would rather say i define equilibrium within the time scale of my experiment
as long as i am doing experiment on the system the thermodynamic variables pressure temperature which i measure they do not depend on time to maintain the equilibrium whatever change i will do is a quasi static change quasi static change means its a very very slow change which means its lower than all other time scales of the problem i can assume that every instant of time my system is in equilibrium i can describe if i have an ideal gas rather one mole of ideal gas i can describe it through the equation of state pV is equal to nRt at every instant of time that explains what i mean
when i plot a pV diagram if i plot p as a function of V i get a pV diagram which says at every instant of time i am in equilibrium i can write pV is equal to nRt okay so quasi static process is very very important to maintain equilibrium and since i am maintaining equilibrium at every instant of time i can write the equation of state ok this is brief recapitulation what i did the concepts of thermodynamics i repeat its a macroscopic subject but eventually the results what we get at the end of the day will be same as those which we obtained from kinetic theory of gas ok let us proceed
so i proposed first law of thermodynamics for you is a conservation of energy
so i can have a situation where i have my system and i give some energy to the system in the form of it
so i provide an energy which i will call thermal energy amount Δq or i use this notation Δq whenever i write tilted delta or q delta it represents an infinite decimal change let me call its a finite change but let me not bother about this nitty gritty details
we use these two notations almost interchangeably ok
so i supply Δq amount of heat to the system and if i do not allow the system to do any mechanical work some quantity must increase and that we call the internal energy internal energy must increase because energy cannot be dissipated
if energy cannot be dissipated some quantity must increase and that is the internal energy let us say we have ideal gas and it is the internal energy of the ideal gas that will actually increase if i provide Δq similarly if i have a system and i push this wall of the container i am doing some mechanical work on the system i am doing some mechanical work i will always assume there is no friction
there is no dissipation that means all i am doing is conservative ok ok if i

do some work again there will be
increase in internal energy because i am not allowing any heat exchange these
two are
the extreme situations which tells me that i have to bother about three
quantities one i
have already written Δq that is heat exchange and then mechanical work
 Δw and i mentioned
something something like a ghost that appears in our theory now that is the
internal energy which i
write as Δu ok which i write as $d u$
so i have Δq Δw and Δu the first
law of thermodynamics which is nothing but the conservation of total energy
will involve these three quantities this should be very very clear
so i must take
into account the internal energy will restrict to two hypothetical extreme
situations in which one
of them i have only heat exchange in the other i only have mechanical work ok
having said this we
proposed the first law of thermodynamics Δq is the heat supplied to the
system Δw is work
done by the system
so Δq should be equal to Δu plus Δw
so whatever heat i supply
to the system will be dissipated in two forms or will be used in two forms one
is system will do
some work
so this is the work done by the system work done as i have written already
let
me write it again work done by the system and this is the increase in internal
energy
so this is how i propose my first law of
thermodynamics all these three quantities Δq Δw and Δu they
should be
put together to get a conservation of energy now if you look at this equation
first
extreme case was no Δw
so case one which we studied in the previous
slide there is no Δw
so you see Δq is equal to Δu
that is what i was trying to emphasize that if there is no work done by the
system
this whole heat i have supplied to the system goes into increasing the
internal energy of the system on the other hand case two which i will soon
show what is that case if Δq is equal to 0 Δu is minus Δw which
tells me that if
system does some work then Δw is positive okay then Δu is negative
because this quantity is positive if this quantity is positive Δu is
negative
which means the internal energy of the system must go down system is doing
some work at the cost of
its internal energy on the other hand if i do some work on this system then
 Δw becomes negative
because of this negative sign Δu is positive that means i am doing some
work on the system and

its internal energy goes up
so this is consistent with whatever we have discussed
so far this
is the first law of thermodynamics and then i briefly mentioned what is an
internal energy
ok internal energy if i consider ideal gas internal energy is proportional to
temperature ok
these i have not proved but for an ideal gas its function of temperature
proportional to
temperature if i increase temperature internal energy goes up this constant
has no relevance
kinetic theory has already taught us that average kinetic energy of the
molecule is related to the
temperature
so internal energy if you consider mono atomic ideal gas molecules then this
quantity
 c_v takes care of counting the degrees of freedom and we already know in
kinetic theory average
kinetic energy is related to temperature
so if you talk about ideal gas molecules your internal
energy is simply translational if it is monoatomic it is translational
vibrational and rotational
if it is diatomic or polyatomic ok
so it is the internal energy means the energy i was talking
about in kinetic theory if i consider ideal gas molecules who takes into care
the fact that
it could be mono atomic it could be polyatomic that is in c_v we calculated c_v
elaborately using
equipartition and that equipartition tells us that whether the molecule is
mono atomic diatomic
or poly atomic goes into calculating the number of degrees of freedom and that
is reflected
in the specific heat capacity
so this is the internal energy and second quantity is it is
extensive that means if you double the size double the number of molecules
definitely it
will get doubled
so it is an extensive quantity as i said in the beginning
so internal energy is
extensive quantity secondly i want to bring a very important notion here its a
state function what
do we mean by state function state function means if i go from a state which
is
characterized by p_i, v_i and t_i to a state p_f, v_f, t_f by means of any
thermodynamic
process ok i will define the thermodynamic processes for you in couple of
minutes this
internal energy only depends upon the initial and final state thermodynamic
variables it
depends upon the values of thermodynamic variables which is initial and final
ok it does not depend
upon the thermodynamic process which is involved in taking the system from
this state to
that state which is very very important

so it is a state function means it is determined by the thermodynamic variables of a state ok for example in ideal gas if you see its proportional to temperature if i make a process in which my temperature goes from t_i to t_f so change in internal energy will be simply $c_v (t_f - t_i)$ it does not depend how i have achieved this process in going from t_i to t_f ok this is very important it is a state function but q and w not so heat absorbed or work done by the system they are not state functions they are functions of the process ok so as i have written here Δq Δw and that is why this Δq Δw depends on the path thermodynamic process ok soon will explicitly calculate them and show that they indeed depend on the thermodynamic process Δu doesn't ok it only depends on the initial and final set as i have calculated here so this is a very important notion and there is a connection between mechanics okay we have already seen adiabatic process this is how thermodynamics brings in the concept of state function or internal energy if you have an adiabatic process ok so in an adiabatic process you know Δq is equal to zero so there is a work done by the system or i am doing some work on the system if i say that there is no friction there is no dissipation what do i have i have a work which is conservative which is conservative this is very very important so i am doing a work which is conservative and now let us recall our mechanics course in mechanics we already know in a conservative force field work done does not depend upon the path i take for example if particle in a gravitational field was initially here i take it here i can take it in many possible ways as many ways you like someone can take it just vertically upward ok but the work done in all the processes are same does not depend on the path rather it depends upon the difference of one quantity in the final state and the initial state that quantity i called potential in our mechanics course ok so potential work done in a conservative force field is in a conservative force field work done is just the difference in potential difference in potential please note that what i am referring as potential is nothing but the so called potential energy or rather i am using the term potential to denote potential energy it does not depend on the path and if i do a work in a closed loop in a conservative force field i start from here and i come back to this point net work done is zero because i come back to

this same position
 potential is the same
 so change in potential is zero that is why my work done will be always
 be equal to zero
 so this is from mechanics and you see in an adiabatic process work done
 is path independent it does not depend on the path because it is a
 conservative force field
 and that gives in the idea of internal energy
 so if you remember if you understand how do we get
 the concept of potential in classical mechanics you immediately get what is
 internal energy
 in thermodynamics ok
 so it is a state function ideal gas as i have mentioned already du
 depends on initial and final temperature because it is a function of
 temperature only
 so if you understand what is the conservative force field and how do we get
 the concept
 of potential in a conservative force field you immediately get to realize what
 is
 internal energy in a thermodynamic process ok and for an adiabatic process it
 is clearly
 a state function if you consider ideal gas because i told you in an ideal gas
 its just $c_v T$ plus some constant this constant has no relevance in essence we
 are
 interested in the difference of internal energy as in mechanics we are
 interested in the
 difference in potential
 so this is the first law of thermodynamics for us this is the
 meaning of internal energy i talked about two extreme processes one is
 completely adiabatic
 only mechanical work is being done other is process which is diathermic in
 which i allow
 heat exchange but general process involve both one is heat exchange other is
 the mechanical work ok now this is the question since i
 have been talking about work
 so much what is the work done by a gas well i will always
 assume quasi static process that means at any step at any instant of time i
 can use my ideal gas
 equation i will always be considering ideal gas not van der waal gas at any
 point on time so
 far if not otherwise mentioned its ideal gas could be one mole could be n mole
 but i will
 always assume quasi static process
 so that at any instant i can use the equation $p v$ is equal to
 $n R T$
 so what is the work done let us assume that there is a container there is
 some pressure some
 p and there is a small shift dx this is the area of the cross section and
 there
 is a pressure p let us say ok now what is the work done work done i know
 force into displacement
 so i am just looking at the magnitude of the work done
 so work done
 should be force which is pressure times area which is dimensionally consistent

and dx is
 this displacement of the wall of the container let us say for example for
 simplicity we are
 taking a rectangular wall okay
 so it is displaced by an amount dx this is the work done which i
 can write as $p dv$
 so i have expressed everything in terms of pressure and volume one intensive
 and
 one extensive thermodynamic variable
 so what will be the network done network done is $\int p$
 dv from v_1 to v_2
 so i am assuming that initially i was as a volume v_1 then the volume
 of the system increases or decreases to v_2 ok i am doing this change
 quasistatically
 so this
 is the work done now if i recall my $p-v$ diagram this is $p-v$ diagram let us say
 so what
 does it imply i am integrating $p dv$ this is my initial volume v_1 this is
 my
 final volume v_2 and this is the work done
 so work done it is simply the area under the curve
 in this $p-v$ diagram ok
 so we have to calculate this area under the curve for
 different thermodynamic processes
 so this is our next time and this will take
 some time and we will try to understand the physical meaning of every process
 so we will
 do it for different thermodynamic processes ok so let us consider for simplest
 thing isothermal
 process what do you mean by an isothermal process temperature is fixed t fixed
 immediately
 you already know one thing if temperature fixed i am talking about ideal gas du
 u is equal to
 zero there is no change in internal energy because in ideal gas internal
 energy is
 just the translational kinetic energy if i consider mono atomic ideal gas
 molecules
 its just translational and its proportional to temperature is temperature is
 kept fixed du
 must be equal to zero now you want to calculate this quantity and you know
 since its a quasi
 static process every instant in time you can write $p v$ is equal to $n r t$
 so that you can write
 this integral as i have written here clearly as $n r t \int \frac{1}{v} dv$ by v ok all
 i have done i
 have substituted for $p v$ here $p v$ is equal to $n r t$ which gives me p is equal
 to $\frac{n r t}{v}$ ok if i substitute it back in this equation i simply get this form and now i
 can integrate
 from v_1 to v_2 this gives me this expression
 so in an isothermal process this is the work
 done there is no change in internal energy so whatever heat is supplied to the
 gas will be
 converted to work done if i give Δq amount of it system will do this much

of work because there
is no change in internal energy
so entire heat supplied whatever i have supplied to the ideal
gas system will go into work done by the system
so this is an isothermal process
important is du is always zero i am considering ideal gas molecules ok
quickly i will go to the next process isobaric again what does it mean
isobaric means pressure is constant pressure is constant which means i will go
from a state $p_1 v_1 t_1$ to a state $p_2 v_2 t_2$ ok its my ideal gas system ok
initial pressure and
final pressure is the same there will be a change in temperature and volume
that is what i need to
calculate and this expression will be different and du internal energy change
in internal
energy is there in this process i cannot say du is equal to zero rather du
not equal to zero
in this process
so in an isobaric process du is not equal to zero remind you all the time
that this Δq and Δw depend on path that means they depend on the
thermodynamic processes
i will always write du to remind you again that u is path independent depends
on the initial
and final state value of u depends on either initial state or final state du
depends on the
difference in initial and final step difference in thermodynamic variables
so isobaric process
pressure is constant we can do this integration very easily ok take the
pressure out of the
integral its a constant
so its $p v^2 - v^1$ its simply p times $v^2 - v^1$ and since
i am talking about ideal gas in both the cases its $p v^1$ is equal to $n r t_1$
in the initial
state $p v^2$ is equal to $n r t_2$ in the final case
so it is simply given by the difference
in temperature
so i have kept pressure constant ok volume and temperature both of these
quantities
change
so p constant v and t change and this is my net work done i can write it
expressing entirely
in terms of change in volume or change in temperature because i am using ideal
gas and i
am doing a quasi static process if i do not do any quasi static process i do a
rapid change i
can go from one state to the other very rapidly but i have to wait initially
it was in
equilibrium with $p_1 v_1 t_1$ finally it may reach an equilibrium which is $p_2 v_2 t_2$
two but i cannot say work done is this because i did not do a quasi static
process rather i
did a rapid change and finally i have to wait for the system to equilibrate ok
this is
another process
so two thermodynamic processes we have already learned one is isothermal

process and second one is isobaric process lets go to the third one isochoric process
so volume is kept constant i am not allowing volume of the container to change immediately
i know that i am not allowing any work no work as i have written here no work is done
neither i am doing any work on the system nor the system is doing any work on the rest
of the universe
so volume is kept constant
so work done must be equal to zero it is
 $p \, dV$ ok work done i again remind you i told you $\int p \, dV$ this dV is zero
if dV is equal to zero no net work done ok
so the heat supplied to the gas is converted to internal energy that is the first example i said that i am supplying some heat
ok if you
remember this represents this situation i was talking about that this situation where i
said that Δq this or this whichever notation you prefer is the amount of heat supplied to
the system but there is no mechanical work ok which means there should be an increase in
internal energy which i am showing here that there is an increase in internal energy recall the first
law recall the first law i was not saying Δq and Δw are conserved ok i am saying take
 Δq Δw and dU if you like all together and then you can have conservation of the total
energy that is my first law of thermodynamics this is why internal energy is essential internal
energy is essential to consume the total energy heat absorbed increases the temperature and hence
the internal energy remind you that i am talking about ideal gas systems internal energy is
directly proportional to the temperature if i supply some heat to the system it increases
the temperature and hence there is an increase in the internal energy these three examples
also establish that work done by the gas depends on the thermodynamic processes
so amount
of work i am doing depends upon the process that is why this notation Δw ok depends
on the path path means whichever process i take the process could include or process could consist of
two or many thermodynamic processes may be first half is isobaric second one is isothermal we will see this type of multiple
processes when i talk about heat engines or obvious
so ok when i talk about heat engine or refrigerators these things will be clear it could be for example isobaric plus isothermal ok multiple
processes in going from one thermodynamic state to the other ok this is an example ok and the work

done will entirely depend on the process i have employed what is important change in internal energy will be the same if i start from the same initial state always and reach the same final state always it just depends upon the initial and final state in precisely speaking depends upon the difference in temperature since i am talking about ideal gas throughout ok now comes the final one what is the work done in an adiabatic process adiabatic process is complicated here Δq is equal to 0 that's important Δq is equal to 0 that means Δu must be equal to minus of Δw this is very important w depends on path ok but Δu doesn't so if i have an adiabatic process then w will be path independent that's what i was trying to say in mechanics conservative force field work done is path independent brought in the concept of potential energy if you like similarly here w is path independent which gives the concept of internal energy ok this is my Δw which is $p \Delta v$ which i have defined now let us recall in kinetic theory we defined c_v and c_p ok ok we used equipartition theorem counted the number of degrees of freedom and then we could easily find out what is the number of degrees of freedom contributing to the specific heat and i used $c_p - c_v$ is equal to R for one mole of ideal gas and i said i will not provide any proof at that point now is the time we can look at it more critically so isochoric process which means volume is constant remember no work done volume is constant if volume is constant i can define c_v that is what calorimetry teaches us that Δq is the amount of heat let us say supplied to the system and Δt is corresponding change in temperature so Δq by Δt keeping v constant this notation means i am keeping v constant let me go this way so this keeps v constant this is $\frac{d u}{d t}$ keeping v constant similarly if you consider a process which is isobaric process in an isobaric process you keep pressure constant ok pressure constant so this is Δq heat supplied to the system Δt increase in temperature but pressure is kept constant c_p this is $\frac{d u}{d t}$ ok what is Δq Δq already we know from first law as i have written there Δu plus Δw which is nothing but Δu plus $p \Delta v$ i am writing Δ here and consistent notation same notation for q u and w but never forget these two depend on path whereas Δu does not so if i use this pressure is kept constant i want to take this derivative pressure constant and limit $d t$ tending to zero what i find out since pressure is constant first

piece will come from internal energy and then pressure is coming out its a constant
 its an isobaric process
 so $\Delta v \Delta t$ holding pressure constant
 so this is my c_v this is my c_p and i want to know what is the difference between c_p and c_v this is important because
 what is our purpose we want to calculate $p dv$ in an adiabatic process in an adiabatic process
 temperature will not remain constant temperature changes it is not an isothermal process if
 temperature changes we immediately know there will be a change in internal energy also
 so change in
 internal energy because of change in temperature and i cannot either use $p v$ is equal to constant
 i cannot say it is not an isothermal process we know $p v$ is equal to $r t$ for an ideal gas one mole
 of ideal gas if temperature is constant then i can say $p v$ is equal to constant ok but it is not an
 isothermal process
 so temperature is not constant ok
 so i cannot say $p v$ is equal to constant we
 have to be careful firstly purpose is to calculate in an adiabatic process t is not constant we cannot use $p v$ is equal to constant which is $r t$ rather for an adiabatic process i
 must have some other equation some other relation relation which is very important
 which i am going to derive for you this relation is $p v^\gamma$ is equal to constant need not be the same constant as the previous one a different constant ok which i will derive for you now for an ideal gas since we are dealing
 with ideal gas we are lucky ok we are lucky for an ideal gas u is a function of temperature
 only
 so this derivative if i write in derivative form its $\frac{d u}{d t}$ keeping volume constant but this
 fellow is a function of temperature only similarly this is $\frac{d u}{d t}$ keeping pressure constant ok so
 these two derivatives since i am assuming internal energy as a function of temperature only
 so these
 two carry no sense ok since internal energy is a function of temperature only these two are same
 now look at this equation what you have c_p what is this quantity you are calculating $p \Delta v \Delta t$
 and you can immediately arrive at the equation this is your c_p this is your c_v all you are left
 with $c_p - c_v$ is equal to $p \Delta v \Delta t$ p and if you calculate it for one mole
 of ideal gas for which you already know $p v$ is equal to $r t$ $p v$ is equal to $r t$ if you
 calculate this quantity this immediately gives you $c_p - c_v$ is equal to r
 this is a

very simple but very insightful proof of this relation that for one mole of ideal gas $c_p - c_v$ is always equal to R that can be proved I have shown in a very simplistic way I repeat the proof c_v is change in internal energy when you change temperature by amount ΔT keeping V constant c_p on the other hand has P constant so it will have a piece this and this space this this is the extra piece now I am arguing that for an ideal gas OK I should always have this quantity equal to this quantity this P and V which I have written here they do not carry a sense because for an ideal gas U is a function of temperature only if that be the case this first term $c_p - c_v$ if you do this term cancels with this term all you are left to it with this term and if you calculate using $PV = RT$ here you simply get $c_p - c_v$ is equal to R so first thing I have proved for you with some basic arguments say $c_p - c_v$ is equal to R now I have to establish this for you I have not defined what is γ for you so I have to define what is γ OK and arrive at this relation so let's proceed so purpose is to know what is constant for an adiabatic process as we know in an isothermal process PV is constant so let us proceed so ΔQ is zero because I am talking about an adiabatic process dU is this quantity which can be easily found out OK this $P \Delta V$ now dU I am talking about one mole of ideal gas so dU is equal to $c_v dT$ that is what I have replaced here you can easily see remember $c_v T$ plus constant that is what I wrote as the expression from internal energy this constant never matters its never relevant because I will always talk in terms of difference in internal energy so this is important any thermodynamic extensive variables like internal energy free energy etc we can always add a constant to that you know in classical mechanics results do not change on where you set the zero of your potential energy so $c_v dT$ is minus $P dV$ this is simply from this equation once you said ΔQ is equal to zero it follows from the first law so I can write in differential form now $c_v dT$ is equal to minus $R T \frac{dV}{V}$ by V what I have done I have substituted P is equal to $R T$ by V let me write probably here P is equal to $R T$ by V OK so I have $c_v dT + P dV$ substituted for $P R T$ over V and already in the previous slide we have proved

for ourselves that $c_p - c_v$ is equal to r
 so $c_p - c_v$ is equal to r and you get an
 expression like this i write it here again
 $c_v \frac{dT}{T}$ is equal to $-\frac{c_p - c_v}{v} \frac{dv}{v}$ times T
 and let us write also Δv in differential form to be consistent okay
 so this is the equation $c_v \frac{dT}{T} = -\frac{c_p - c_v}{v} \frac{dv}{v}$
 $\frac{dT}{T}$ is equal to $-\frac{c_p - c_v}{v} \frac{dv}{v}$ into $\frac{dv}{v}$ ok everything is
 in differential form
 now and now i can finally arrive at this equation now i have intentionally
 not
 written where is γ here you can immediately see γ is nothing but
 $\frac{c_p}{c_v}$
 so ratio of specific heat capacity at a constant pressure to the specific
 heat
 capacity at a constant volume that is your γ that is $\frac{c_p}{c_v}$ well
 so this is the equation
 you have now i write it again for clarity $\frac{dT}{T} = -\frac{1 - \gamma}{\gamma} \frac{dv}{v}$
 you can integrate if you integrate you get this result its very simple
 integration you always
 know $\frac{dT}{T} = \frac{1}{T} dT$ if you integrate you get a $\log T$ ok thats what precisely i have
 written here $\log T$
 $\log T - \frac{1 - \gamma}{\gamma} \log v$ plus some constant which comes from integration and
 so that gives me T is
 proportional to v to the power $1 - \gamma$
 so i have a relation in an adiabatic process ok
 i can put it in more succinct form if i proceed i had this relationship which
 was telling me T is
 proportional to v to the power $1 - \gamma$ T is proportional to v to the
 power $1 - \gamma$
 or i can write T is equal to some constant $\frac{v}{T}$ to the power $1 - \gamma$
 now use $p v$ is equal
 to $r T$ because i am using one mole of ideal gas one mole of ideal gas i use $p v$
 is equal to $r T$
 this relation is $p v^\gamma$ is equal to constant ok the question i asked
 question i started with for an isothermal process
 you have $p v^\gamma$ is constant it is a quasi static process at every instant
 of time its in equilibrium and its your $p v$ is equal to constant this is
 your isothermal process what is the corresponding relation in an adiabatic
 process
 remember adiabatic process is also quasi static but temperature is not constant
 rather you will
 have $p v^\gamma$ is equal to constant ok for an adiabatic process $p v^\gamma$ is
 equal to constant
 that is what we have to use when we want to do this work done in an adiabatic
 process
 ok remember this constant if i called c isothermal this constant is c adiabatic
 ok
 maybe i can write this in the following way that in an isothermal process $p_i v_i$
 it does not refer to initial state rather it refers to a process ok maybe
 to be precise i use capital C
 so $p_i v_i = C$ that means isothermal process is
 a constant C_i for an adiabatic process $p v^\gamma = C$ adiabatic process on the other hand you

have

p v γ some other constant c ok the subscripts here refer to the process not the initial value now if i give you two isotherms or two p v diagrams

not isotherms i correct myself adiabatic process temperature is not constant so i give

you two curves two p v diagrams and ask you the question if one is adiabatic and

other is isothermal which one is isothermal which one is adiabatic i leave this question

with you and in the beginning of the next lecture we will come to that that given two curves in

the p v diagram on the p v plane ok i am telling you one is isothermal and one is adiabatic you

have to tell me which one is adiabatic and which one is isothermal but by looking at the curves or

rather looking at the slope of the curves this is a hint

so our initial purpose was to calculate the

work done in an adiabatic process and adiabatic process is complicated you are going from p one v

one t one to p two v two t two ok earlier remember earlier we were keeping one of them fixed

in isothermal process i was having t fixed isochoric process i was having v fixed ok in

isobaric process i kept p fixed but when i have an adiabatic process all the thermodynamic variables

change but there is an important simplification p one v one γ should be equal to p two v

two γ that is what i have proved for you well now what is the work done i have to

calculate w v one to v two and then p d v which now i can write since

this is always satisfied it is a quasi static process every instant this is satisfied i can write it as c adiabatic and then d v by v to the power γ v

one to v two ok now

all you have to do you have to do this integral and if you do this integral you will find out work done i am not doing the integral for you

you can very easily do that one minus γ p two v two minus p one v one this is the

work done in an adiabatic process automatic process is very complicated because all the

thermodynamics variable change but keeping this constant and that enables us to find out a

closed form expression for the work done ok and this is the expression for the work done and now

p v is equal to rt is always satisfied after all it is ideal gas

so it is satisfied initially

also in finally ok

so you can write this as r t one minus t two by γ minus one ok this is the final answer and γ is always

greater than one because c p exceeds c v

so this is the work done in an adiabatic process and i

will come back to this in the next lecture and what is the essence of today's lecture

i talked about work done by the gas and established for you that it depends

upon

the thermodynamic processes ok whichever process you take it depends on the thermodynamic processes

for example it is 0 in an isochoric process most complicated is the adiabatic process i have given you an expression in the next lecture we will start again from

the adiabatic process thank you for today you

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