

welcome to the fifth lecture of the series of lectures on kinetic theory and thermodynamics this lecture hour will be essentially spent on describing the basics of thermodynamics but as our regular practice is i will be doing a little bit of recapitulation of kinetic theory of gas what we did in the last lecture we talked about mean free path and non ideal gas both the things i will touch upon briefly before i move on to thermodynamics so we calculated the mean free path what is this quantity this quantity is the distance traverse by a molecule between two successive collisions since we are talking about kinetic theory i need not remind you this is the average distance and this quantity is actually given in terms of average velocity and time between two successive collisions which we calculated and that gives me the mean free path which is $\frac{1}{n \pi d^2}$ where d is the diameter of the molecule and n is the number density now here we assume that molecules are hard spheres with a diameter d then what did we do we constructed a cylinder this is the height and this area is πd^2 and molecular diameter is this d and this radius is $\frac{d}{2}$ ok now we draw this picture this is the top view of the cylinder if you like this is the molecular radius $\frac{d}{2}$ and this is the radius $\frac{d}{2}$ of the cylinder i have constructed so what is important is whenever any other molecule let me remind you that i assumed other molecules are static in the beginning whenever the center of any other static molecule comes here tries to penetrate these cylinder there is a collision so whenever any other static molecule has is centered either on this line or inside that is inside the bigger cylinder there is a collision using this we could calculate the total number of collisions in time Δt total number of collisions we found out that it is $n \pi d^2 \Delta t$ this is total number of collisions over a time Δt from that one can easily find out what is the time between two successive collisions and hence the mean free path which has this form $\frac{1}{n \pi d^2}$ so it is important that molecule has a finite size and that is taken into consideration while calculating this mean free path now there could be many questions here this target molecule which i mentioned in the last lecture it suffers collisions though it gets deflected right it should be deflected because its having collisions if it has collisions then can i talk about a

single cylindrical geometry which i am mentioning here it is not actually the fact of course it gets deflected but on an average sense locally i can assume still there is a cylindrical geometry and whichever static molecule does have its center lying within the bigger cylinder will suffer a collision

so on an average i can assume there is a cylindrical geometry of area πd^2 and height $v \Delta t$

so this gives me the expression for mean free path there is an additional approximation that approximation is other molecules are static which is never the case one should not take v average rather one should be taking relative velocity between two molecules under consideration if you do things in a more rigorous way you will find out a correction of factor of root two coming into the expression of mean free path

so i wanted to emphasize the following point that we are making an assumption that on an average there is a cylindrical geometry even if my target molecule is suffering many many collisions now then we went to non ideal gas non ideal gas is a van der waal gas for van der waal gas the equation of state for one mole of van der waal gas is given like this ok so this is very important that you note that there are two corrections one correction is a by v^2 other correction is b okay first correction as i told you is a correction to pressure and this fellow a by v^2 must be having dimension of pressure where does this correction come from i told you that i am assuming molecules are completely non interacting but remember there is a weak attractive interaction between the molecules this a by v^2 captures that it is actually the average of these attractive interaction i gave more specific example by saying if a molecule hits the wall is at the wall this molecule will be pulled by an attractive force by all the other molecules which are inside the container and number of molecules which are on the wall at any instant will be proportional to n by v similarly number of molecules which are inside the container will also be proportional to n over v

so average correction to pressure will be of the form constant by v^2

so this correction arises due to attractive nature of interaction which is very weak and when a molecule is inside the container i can conclude that on an average this is zero but when it is at the wall there will be an attractive force which effectively corrects pressure now second correction was b why does

become

b comes because of the fact that molecules as i also had explained in the case of mean free

path that molecules are of finite size we assumed they are actually hard spheres of diameter

d in that case i explicitly calculated with some phenomenological arguments that b should

be proportional to 4 times number of molecules in the container and then volume of a particular

molecule which i calculated i showed assuming the spherical molecules this expression is

$\frac{4}{3} \pi d^3$ this is the volume of a molecule ok

so b is proportional to volume

of the molecule what does b represented for us b tells us if i take a molecule the entire volume

is not accessible to it each molecule has a volume which is excluded i say it if this is a molecule

and took two molecules i can assume a spherical concentric spherical volume which has radius d

which is excluded for the other molecules ok so excluded volume correction is incorporated in this

parameter b and i have my van der waal equation of state which is p plus a by v^2 square correction

to pressure due to average attractive interaction b correction to volume because the entire

volume is not accessible to the molecule each molecule has a finite size hence a

finite amount of volume will be excluded for any molecule i have in consideration and then

this b will be naturally proportional to number of molecules i have in the container and volume

of each molecule ok that explains the origin of a and b now if someone wants to generalize to n

moles ideal gas we know $p v$ is equal to $n r t$

so what should be the form in the case of van der

waals gas or the real gas let us assume n moles ok for n moles you know correction to volume will

be more because i have more molecules now in the same volume

so my excluded volume will be n times

more excluded volume will be n times more ok that is what is reflected v minus $n b$ and then here

also density goes down by a factor of n

so its n and v into n by $n v$

so you will be having

a n^2 square term coming here

so your van der waal equation i write it down again will

now turn out to be v minus $n b$ is equal to $n r t$ then the limit a going to 0 that means when

you ignore you are allowed to completely ignore the interaction this attractive force a is equal

to zero when you can assume the point particle nature of the gas molecules then you can

ignore b and when you set a is equal to zero is equal to b you retrieve your ideal gas

equation ok but these two corrections are very nontrivial corrections ok and these allows us to explain the phase transition for example the liquid gas phase transition that we often talk about and that we encounter in our real life okay so having said that i wanted to bridge a connection between the van der waal equation of state and the phase transition liquid gas phase transition ok so i drew the isotherms water isotherms i fix the temperature ok i fix temperature plot pressure as a function of volume and i claimed these are the curves well i wanted to note something here which i not intentionally mention in the last lecture if you really plot van der waal equations mathematically you can use it in your computer you will see the curves are not really like that typically a isotherm will look like this ok without much explanation which is far beyond your syllabus i will say there is something called maxwell's construction which gives this form which i have drawn and this form is experimentally verified what i said here there are different temperatures t_1 t_2 t_c let us say and this is term t_3 which is greater than t_c temperature is increasing so t_2 is greater than t_1 t_c is greater than t_2 and so on so there is a liquid phase which is high pressure low volume there is a gaseous phase which is high volume low pressure so i can go from one phase to the other by changing temperature or by changing pressure but there is a difference there is this temperature t_c and i draw drew a dotted line here i said this region is actually the coexistence region coexistence region in which liquid and gas coexist i have to correct myself liquid and vapor coexist why it is important because i gave you the definition of vapor vapor is the gas below this critical temperature t_c as i repeatedly mentioned above the critical temperature t_c which is this t_3 remember i am plotting isotherms p as a function of v for different temperature so if i look at the isotherm at the temperature t_3 which is higher than t_c no amount of pressure can liquefy the gas its always in the gaseous state so liquefaction of a gas by application of pressure is not possible if i am at a temperature which exceeds the critical temperature so van der waal equation two simple corrections we gave two simple corrections one to pressure coming due to the fact of attractive interactions finite size gives one more

correction that is explained the origin of these two corrections

these correction terms have drastic effect on the physics i am studying it explains liquid gas transition you see through a coexistence region ok below T_c there is a coexistence region i change pressure i go from liquid phase to the gaseous phase or the vapor phase through a coexistence region in which liquid and vapor coexist and above the critical temperature no amount of pressure can liquefy a gas ok if you like one can draw a p versus T diagram keeping volume fixed you can see it coexists on this coexistence line that ends at this critical temperature

okay this also i explained so phase transition can only come through interactions and van der waal equation simplest possible correction to ideal gas equation gives us phase transition which is very very important having recapitulated whatever i said in the last lecture i will be now moving on to describe thermodynamics and basics of thermodynamics

there is a fundamental difference between kinetic theory and thermodynamics which i explained in the very first lecture thermodynamics is a macroscopic approach to understand the properties of matter macroscopic approach what do we mean by that i do not care about what is happening at the molecular level i do not care about the velocity distribution or speed distribution or average velocity all i will look at is experimental measurable quantity ok i will be looking at only pressure volume temperature which i measure experimentally so its a coarse grain description in that sense which i mentioned that means i am looking at the macroscopically measurable objects like pressure volume temperature remember in kinetic theory also we are talking about average speed average velocity but eventually everything was connected to pressure volume and temperature

so kinetic theory had its own definition of temperature that is given in terms of average kinetic energy of the molecules similarly thermodynamics will have its own definition of temperature but it will be again connected to experimentally observed temperature in absolute or kelvin scale which i introduced in the very beginning of this set of lectures so its a coarse grained description does not probe the system at the molecular level we do not care what is happening at the molecular level and thats why i call it a macroscopic approach so v T these i will call thermodynamic variables well you know degrees of

freedom in mechanics you learn degrees of freedom position and momenta when i discussed equipartition theorem i talked about x and p both contributing half kT to the energy here in thermodynamics there is no x no p all degrees of freedom i will be talking about are p v p and there are other quantities for other systems but we will restrict ourselves to p v and T ok now basic idea is this thermodynamic approach talks about a system a system plus the rest of the universe so i will have a system which my experimental system and then the rest of the universe which is interacting with the system ok this is very important statement so i will have a system this is my S and then i will have the rest of the universe which i will call the resolver ok so i have a system and a resolver and they are separated by a wall ok so you see the difference i do not know what is there what molecules are doing but i know there is a system which is described by pressure volume and temperature may be chemical potential let us not go into that and this is separated from the rest of the universe whatever else is there in the universe with some wall and these walls play very important role in defining thermodynamic variables which i will come soon the walls determine what type of interaction the system will be having with the rest of the universe which i will simply call universe subsequently and the equilibrium that system reaches now since i mentioned thermodynamic variables i must say that they are of two kinds one is extensive other is intensive let me define what do i mean by extensive and intensive thermodynamic variables let us take a big system described by n v T p and things are in equilibrium which means that nothing depends on time so they have a fixed value constant in time value now if i divide it into two halves is the volume v by two v by two ok now volume is halved what happens to number of particles i am assuming the system is divided into two equal halves in equilibrium so whatever equilibrium was there before i am making this division that equilibrium is maintained so if pressure was p here that let me see this is the container one of the wall of the container this pressure was p for equilibrium ok so pressure will remain the same when i divide into two halves temperature again if i have to keep things in equilibrium that is nothing depends on time temperature should not change but you see v goes to v by two similarly n goes to n by two so you see that there are some quantities which become half of their initial value there are some quantities which are not

at all influenced or affected by this division of the system into two halves
 ok
 quantities which are indicative of system size that means if i have the system
 they get halved they become half of their initial value these are called extensive
 quantities for example volume number of particles and also a quantity which i am going to define very
 soon is called the internal energy these are extensive quantities ok for example if i define
 external energy it is a function of volume sorry this is the internal energy the
 internal energy which is a function of volume if i increase the volume by a
 factor \times three times more
 so volume let us say volume goes to \times volume \times can be any number it could be
 two three half in the present example internal energy actually will show that goes
 to \times
 u these are extensive quantities maintaining equilibrium you make the
 volumes \times times the initial value the internal energy number of particles they
 all go \times times internal energy u goes to $\times u$ let me be more precise i am saying
 increase the volume of the system maintaining equilibrium the simplest way to think that i am
 considering a bigger volume of this same system okay as i have shown in this example henceforth whenever
 in the context of extensivity i say increasing the volume maintaining equilibrium i mean i am
 considering a bigger volume of the same system question is what are other types of variables
 these are intensive for example pressure temperature they do not have any change because of this
 multiplying by a factor of \times
 so these are intensive quantities so intensive quantities are completely
 insensitive to the size of the system whereas extensive quantities are indicative of the
 size of the system maintaining equilibrium if i double the system all these extensive
 quantities will become double but there is an interesting thing density what is density density is n
 by v ok if volume goes to $\times v$ number of particle which is also an
 extensive quantity will go to $\times n$ so \times cancels with \times you see density remains
 the same
 so density is an intensive quantity true for any ratio of two extensive
 quantities whenever i take the ratio of two extensive variables that becomes an intensive quantity
 ok
 so this was important to tell you that thermodynamic variables they describe entire
 thermodynamics and they are of two type one intensive and other is extensive having said this lets go to
 the walls i

told you i will be having a thermodynamic system and this thermodynamic system will be separated from the rest of the universe or simply the universe as i will say and this is my system separated by a wall ok this walls will determine what type of interaction exists between this system and the universe first is adiabatic world what do you mean by an adiabatic wall let us come to this adiabatic wall means this wall is completely non-conducting so system is insulated from the rest of the universe system is insulated from the rest of the universe this implies there is no heat exchange here no heat exchange this is called an adiabatic wall ok this is very important there is no heat exchange then how can this system interact with the universe it can interact with the universe only through mechanical interaction i can have this wall movable if i move this wall then there is some energy supplied to the system i am doing some work on the system ok thats how it can interact with the rest of the universe ok then there is a diathermic wall diathermic wall is just opposite of what i defined as an adiabatic wall diathermic wall on the other hand allows for heat exchange this is very very important one wall is not allowing any heat exchange system is completely insulated that is adiabatic wall system is completely insulated from the rest of the universe ok diathermic all on the other hand it allows for heat exchange ok so there is mechanical interaction in the case of adiabatic world which is possible and then in the diathermic wall there is a thermal interaction possible that means there is heat exchange also mechanical interaction is not stopped so we can in general have a wall in which you will have both thermal interaction and mechanical interaction but these are two idealized situation in one there is no heat exchange possible in the other heat exchange is possible also it is not part of our syllabus will not go into that there could be porous walls which allow for particle exchange ok you can allow particle exchange and then reach some situation when chemical potential of the system and the universe become equal and then there is a equally an equilibrium reached ok but i will not discuss anything about this chemical interaction we will restrict ourselves to adiabatic walls and diathermic walls ok and talk about thermal interaction means heat exchange and mechanical interaction that is i am moving this wall of the container ok so let us proceed what is equilibrium ok i have told you that

there are walls and walls give my interaction now question is what is equilibrium equilibrium
i defined in the beginning of the kinetic theory lecture also that nothing depends on life i
measure measure pressure at time t is equal to $t = 0$ if it is p and then i measure at t is equal to
 $t = 0$ even then pressure will be p it does not depend on time ok
equilibrium is an idealized
concept but we will always assume that system is in equilibrium and nothing depends on time
so i
define equilibrium constant value of thermodynamic variables constant in time they do not change
with time i told you thermal interaction then that gives me notion of one equilibrium
that is called thermal equilibrium what is thermal equilibrium we already know that there is
a quantity called temperature
so i have a system and the rest of the universe they can exchange heat if it is a diathermic wall ok and when they reach the equilibrium we
know basic notion of temperature tells us in equilibrium this temperature of the system should be equal to the temperature of the rest of the universe
so T_s is equal to T_u ok there is no heat exchange there is no energy exchange equilibrium has reached and since reservoir wire is very big i can say i give the term reservoir for you its a very big thing rest of the universe
ok
and i can say it has infinite heat capacity its also an idealized concept but very useful infinite heat capacity if it has infinite heat capacity its temperature
does not change
so there is exchange of heat between the universe and the system when equilibrium is reached temperature of the system will be equal to the temperature of the reservoir
that is the temperature at equilibrium
so that is thermal equilibrium
so thermal equilibrium means
that temperature is equal between the system and there is a wire there is no further heat
exchange what is then let us say mechanical or adiabatic situation if you like
so in
the mechanical equilibrium what will happen i will show in more details later that pressure
ok pressure i can move this wall of the container there is a pressure of the system
there is a pressure of the universe i can move this container in such a way ok that in equilibrium p_s is equal to p of the universe
so pressure should be equal
remember i am always equating some intensive variables in one case thermal interaction it was
the temperature other case mechanical interaction it is the pressure which becomes equal and i
say there is a mechanical equilibrium if you like achieved ok but most of the

cases i will be talking about both mechanical and thermal that means i will be allowing heat exchange as well as mechanical interactions ok and then system reaches an equilibrium and i will do the thermodynamics of that system which is now in equilibrium

so let me summarize what i said there are walls walls are separating the system from the universe the walls can be diathermic or adiabatic in the adiabatic situation there is no heat exchange diathermic situation there is heat exchange once determined also what type of equilibrium i will have for example if i allow heat exchange then equilibrium will reach when the temperature of the universe will be equal to the temperature of the system that is equilibrium situation no further heat exchange and i work can work on the with the system which has reached the equilibrium similarly one can talk about a mechanical equilibrium in which the wall of the container is movable i move it in such a way that pressure is balanced pressure is same between the system and the universe that is the mechanical equilibrium ok and i will be talking about situation which has reached mechanical thermal equilibrium nothing depends on time and then i will do the thermodynamics of that particular system ok now having said all this walls mechanical interaction heat exchange i will establish for you rather i will propose the first law of thermodynamics and i as i said first law of thermodynamics is nothing but the conservation of energy if you remember these i mentioned maybe in my second lecture

so first let consider let me consider a container inside that you can assume this gas molecules i have been talking about in kinetic theory they are moving around but in thermodynamics i do not want to know their velocity distribution ok now lets say i supply an amount of heat Δq or let me use this notation Δq ok an amount of heat Δq i supply to this system what would happen to this system this energy goes up but if i do not allow any mechanical interaction ok i do not allow any change in volume what will happen to this energy supply to the gas system something should increase that is my conservation of energy tells me energy cannot be dissipated ok

so you have to keep in mind something additional is there other than mechanical energy and thermal energy i have been talking about

so first i can change energy by supplying heat by supplying heat then forget about heat which i already explained you move this

wall you move this wall with some velocity u
 so what will happen this is moved ok you can
 roughly think that molecules coming with some velocity here hitting the wall
 but unlike kinetic theory where i was assuming this wall is a static object this molecules will
 not go back with the same velocity or same speed because there is now a relative velocity
 wall is also moving ok
 so since wall is also moving the molecule which hit this wall will go back
 with a different speed
 so its kinetic energy changes ok this is a very heuristic very
 phenomenological idea but all i wanted to say by doing this mechanical work by moving the
 wall of the container i am changing the energy of the system
 so energy of the system can be changed in two ways one by supplying or extracting heat i can take heat
 away from the system or i can supply some heat to the system energy increases and then this
 mechanical part of the energy which tells me if i move the wall ok if i move the wall what
 will happen that energy will change that i very roughly very approximately said that if molecules
 coming with average speed v they will not go back after hitting the wall with the same average speed
 there will be a change because this fellow the wall itself is moving with a velocity u
 so there is a change in energy
 so change in energy can be achieved also due to mechanical work this
 is what i learnt in mechanics ok i do some work ok
 so i get some change in energy now question is i am changing how what is the process am i changing the wall of
 container very fast very rapidly ok that is not true you i am not really doing
 that all the processes i have talked about heat exchange i have talked about
 mechanical interaction in which i change the wall of the container move the wall of the
 container but i am not doing it very very rapidly this brings in the concept of quasi
 static process ok what is a quasi static process quasi static process means its a
 very very slow process how slow is flow you can ask me it is very slow in the sense you draw
 a pv diagram right you have a pv diagram ok given a p you always give me a value of v i choose a
 value of p i immediately get a volume value of volume v ok and i expect them to satisfy the
 equation for an ideal gas pV is equal to RT and i said this equation is valid only in
 the equilibrium so quasi static process is a process you should understand its quasi is static that
 means almost static i

am changing the parameters be it movement of wall or supply of heat very very slowly there is an infinitesimal change slower than any other characteristics time scales of the problem ok any other characteristic time scale of the problem which is very slow more importantly i am making a change but every instant i can assume that system is in equilibrium so this is a quasi static process a quasi static process means its a very slow process and every instant of time i can assume system is in equilibrium for an ideal gas i can write pV is equal to nRt okay this is very important notion which will keep coming back to ok so whatever changes i mentioned here are all quasi static changes whatever i am saying will completely break down if i do very rapid process ok what happens if i do a rapid process ok there will be final state reached i have to wait for the system to equilibrate when system equilibrates once again all thermodynamic variables reach a time independent value i can do thermodynamics but in between what happen i do not know but what is the difference in the pV diagram at every instant of time i am assuming that system is in equilibrium and i can write pV is equal to nRt so this is a very important notion walls equilibrium and to be always in equilibrium i need quasi static processes whatever process mechanical thermal i am talking about are all quasi static processes ok now i have said that the two types of energy which i will talk about one is thermal exchange exchange of heat i wrote it as Δq let me say the mechanical work this is Δw ok now should i then say Δq is equal to Δw and that is my energy conservation no i will be in trouble because here i am just allowing exchange of heat there is no work done on this system so heat must go to some other form of energy ok this point should be clear here i am not talking about any work no work if no work there is no mechanical energy so there must exist some other form of energy in which the heat is getting converted to similarly here if i do not allow any heat exchange what would happen there must be some other form of energy to which this mechanical energy is going ok remember here if no work is done in this process i am supplying heat some energy of the gas must go up and that energy is called the internal energy ok this is a very important concept that if you do not do any work here system is not allowed to do any work so where does the heat energy go heat energy goes in increasing the

so called internal energy of the system similarly here if you do not allow any heat exchange the mechanical work that you do on the system goes on to increasing the internal energy of the system so this is first fundamental notation or fundamental concept which i brought in which is called the concept of internal energy so when you talk about conservation of energy when you talk about this conservation of energy please remember the conservation of energy includes heat energy thermal energy here mechanical energy here or the work done and the change in internal energy with all this preamble system resource wire walls interactions okay notion of internal energy quasi static processes i now put before you what is the first law of thermodynamics this is the first law of thermodynamics which is written out on this slide you can see Δq heat supplied to the system i am supplying an amount of heat which is Δq and then Δw is the work done by the system see previous example i took two extreme cases here i said heat supplied no work here i said i am doing some mechanical work on the system not allowing any heat exchange here i am doing both that is why in one of the previous slide set i said that i will be talking about both mechanical and thermal interactions and then system reaches the equilibrium i will deal with that particular system and do heat exchange or mechanical interaction in a very quasi static way so that i can always assume that system is in equilibrium so Δq is the heat supplied to the system Δw work done by the system and then first law of thermodynamics Δq is equal to Δw plus this new quantity which is called Δu that i will be calling the internal energy ok so you see if Δw is 0 first example of the previous slide Δq is Δu so whatever heat i supplied went into increasing this internal energy ok and then if i do some work but do not allow ok do not allow any heat exchange so that Δq is zero so you see Δu is equal to minus Δw ok minus Δw and this is where i need to fix a convention ok i will fix the convention in the following way Δq positive when heat is supplied to the system internal energy increases Δw is positive work done by the system ok in the previous example if i was doing some work ok on this system then it would be negative and internal energy goes

up

so i repeat the first law of thermodynamics Δq heat supplied to the system Δw work done by the system ok and then again i write it in the form which is written here that is Δq is equal to Δu plus Δw okay convention is Δq is positive heat supplied to the system heat is extracted from the system Δq will be negative ok if you set Δw is equal to 0 Δq is equal to Δu

so whatever heat you supply to the system it goes into increasing the internal energy of the system if you extract heat from the system then Δq is negative internal energy goes down ok and now other way around if you talk about this one you can see from this expression in one case internal energy will increase when i work on the system ok then this Δw itself will be negative from the equation for the first law of thermodynamics i get this one if Δw is negative means i am working on the system internal energy increases if Δw is positive that means system is working at the cost of its internal energy

so you see for energy conservation internal energy must be there we should understand what is this internal energy ok before i proceed to that let me write it in the following form the first law of thermodynamics what i have written here which is actually conservation of three quantities thermal energy work done or mechanical energy and the internal energy can be written in a differential form ok it says Δq Δw and $d u$ this is where one has to be careful why these two are Δ and this one is d these i will explain briefly today and elaborately in the next lecture but let me proceed a little bit what is the internal energy i said if supply heat to the system and do not let it work its internal energy goes up what is this internal energy ok you know if you supply heat kinetic theory has already told me that average kinetic energy increases which implies that temperature increases ok without further proof at the moment i will just say let us say ideal gas ideal gas molecules let us say monoatomic if it is mono atomic ok i know the translational kinetic energy if i take lets say one mole of monoatomic ideal gas $\frac{3}{2} n k$ b t here it is the avogadro number which we kept referring to so this translational kinetic energy is actually the internal energy how do you know this analogy clearly you know this analogy clearly because you know when you supply heat temperature increases kinetic

theory teaches us that average kinetic energy increases translational kinetic energy ideal gas mono atomic only translational so for mono atomic ideal gas it is the translational kinetic energy that is the internal energy so you increase the temperature what goes up is the translational kinetic energy of the monoatomic cache molecules ok so what is internal energy internal energy is equal to this form which i will not prove probably but i will expect because of this analogy it is $c_v t$ plus constant for ideal gas and you know c_v has the information whether ideal gas is mono atomic diatomic or poly atomic i count the degrees of freedom taking either translational or translational plus rotational or translational plus rotational plus vibrational all those informations go into this c_v so its the kinetic energy if you talk about ideal gas molecules ok this is a good point i think where we should stop but i have to tell you this is a very important thing this Δq Δw and Δu actually Δq and Δw they depend on the thermodynamic processes i have already shown you two examples in which i have in one i have Δq is equal to \emptyset only Δw was there in the other there was a Δq but no Δw so this Δq and Δw if i go from a initial state to final state ok let us say this is p_i this is p_f i go from an initial to final state let me p_i to p_f ok or v_i to v_f this Δq and Δw depend on how i have reached from the initial to the final state but Δu does not depend on how i went from the initial to the final state rather it depends upon only the initial state and the final state ok that is a very important concept which i will explain from the point of view of mechanics where you already know there is a concept of potential for a conservative force field that i will generalize here to explain to you what is this internal energy and what do i mean by the terminology that it is a state function so i end this by saying that this quantity actually is a state function this quantity u is actually a state function depends on initial and final state that means initial and final values of the thermodynamic variables if it is ideal gas i already have tried to explain to you that it is $c_v t$ so it will actually be given by the difference in temperature of these two states ok thank you for today you